# 01 THE IMPLEMENTATION OF A HAIR SCATTERING MODEL

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Although the third edition of *Physically Based Rendering* includes an implementation of a shape that allows for efficient intersections of rays with a "flat ribbon" primitive that can be used for modeling fine hair and fur, the system doesn't include any BSDFs or BSSRDFs that model light scattering from hair. (This state of affairs was due to both time and space limits.)

Therefore, here we will describe the implementation of the hair scattering model described in the paper A Practical and Controllable Hair and Fur Model for Production Path Tracing, by Matt Jen-Yuan Chiang, Benedikt Bitterli, Chuck Tappan, and Brent Burley. The paper itself is available here: benedikt-bitterli.me/pchfm/. See the "Further Reading" section at the end of this document for more information about previous work in scattering from hair (some of which is also incorporated in our implementation here.) Figure 1.1 shows a model of curly hair rendered using this BSDF.

This implementation includes both a new BxDF and a new Material for hair. Both of these are defined in the files materials/hair.h and materials/hair.cpp, which are now included in the "master" branch of pbrt. A few unit tests are in the file tests/hair.cpp. Note that this is a different organization than the rest of pbrt, where BxDFs are generally defined in core/reflection.{h,cpp}. For this extension, we wanted to localize the additions to all-new files as much as possible in order to minimize changes to preexisting parts of the system.



## 1.1 GEOMETRY

Before discussing radiometry and light scattering from hair, we'll start by defining some ways of measuring incident and outgoing directions from intersection points on hair.

We will assume that the hair BSDF is only used with the Curve shape that was defined in Section 3.7. in the third edition of *Physically Based Rendering*.<sup>1</sup> A Curve can represent the shape defined by circle swept along the path of a Bézier curve, giving a generalized cylinder, and provides a reasonably efficient intersection test for this primitive. For the geometric discussion to follow, we'll assume that the Curve variant corresponding to a flat ribbon that is always facing the incident ray is being used. However, in the BSDF model, we'll interpret intersection points as if they were on the surface of the swept cylinder. If there is no interpenetration

<sup>1</sup> All page and section references in the remainder of this document will refer to the third edition.



between hairs and if the hair's width is not much larger than a pixel's width, there's no harm in switching between these interpretations.

Throughout our implementation of this scattering model, we will regularly find it useful to separately consider scattering in the longitudinal plane, effectively using a side view of the curve, and scattering the azimuthal plane, considering it head-on at a particular point along it. To understand these parameterizations, first recall that Curves are parameterized such that the u direction is along the length of the curve and v spans its width. At a given u, all of the possible surface normals of the curve are given by the surface normals of the circular cross-section at that point. All of these normals lie in a plane; we will call this the *normal plane* (Figure 1.3).

We'll find it useful to represent directions at a ray-curve intersection point with respect to coordinates  $(\theta, \phi)$  that are defined with respect to the normal plane at the *u* position where the ray intersected the curve. The angle  $\theta$  is the *longitudinal angle*, which is the offset of the ray with respect to the normal plane (Figure 1.4(a));  $\theta$  ranges from  $-\pi/2$  to  $\pi/2$ , where  $\pi/2$  corresponds to a direction aligned with  $\partial p/\partial u$  and  $-\pi/2$  corresponds to  $-\partial p/\partial u$ . As explained at the start





of Chapter 8, in pbrt's regular BSDF coordinate system,  $\partial p/\partial u$  is aligned with the +x axis, so given a direction in the BSDF coordinate system, we have  $\sin \theta = \omega_x$ , since the normal plane is perpendicular to  $\partial p/\partial u$ .

In the BSDF coordinate system, the normal plane is spanned by the y and z coordinate axes. (y corresponds to  $\partial p/\partial v$  for curves, which is always perpendicular to the cure's  $\partial p/\partial u$ , and z is aligned with the ribbon normal.) The *azimuthal* angle  $\phi$  is found by projecting a direction  $\omega$  into the normal plane and computing its angle with the y axis. It thus ranges from 0 to  $2\pi$ . (See Figure 1.4(b).)

One more measurement with respect to the curve will be useful in the following. Consider incident rays with some direction  $\omega$ : at any given parametric u value, all such rays that intersect the curve can only possibly intersect one half of the circle swept along the curve (Figure 1.5). We will parameterize the circle's diameter with the variable h, where  $h = \pm 1$  corresponds to the ray grazing the edge of the circle, and h = 0 corresponds to hitting it edge-on. Because pbrt parameterizes curves with v across the curve and  $v \in [0, 1]$ , we can compute h = -1 + 2v.



Given the *h* for a ray intersection, we can compute the angle between the surface normal (which is by definition in the normal plane) and the direction  $\omega$ , which we'll denote by  $\gamma$ . (Note: this is unrelated to the  $\gamma_n$  notation used for floating-point error bounds in Section 3.9). See Figure 1.5, which shows that  $\sin \gamma = h$ .

## **1.2 UTILITY ROUTINES**

Before going forward to the hair scattering model implementation, we'll introduce a few utility functions that will be repeatedly useful in the following. (In retrospect, these would have been nice to have included in pbrt-v3; they will likely be part of the core system in a future version.)

First, Sqr() just computes the square of the given value. Although this function provides trivial functionality, it makes it possible to transcribe equations to code more succinctly than if we did not have this helper.

```
\label{eq:General Utility Functions} \\ \equiv \\ \mbox{inline Float Sqr(Float v) { return v * v; } } \\
```

In the following, we will also need to compute relatively large integer powers of floating-point values. Because the integer powers will be compile-time constants, it's possible to use C++ templates to generate much more efficient code to compute these powers than is generally possible using standard library routines.

Consider for example the task of computing the value  $v^{20}$ . This computation is equivalently expressed as

$$(((v^2)^2)^2)^2(v^2)^2$$

Counting up multiplications, we can see that given v, just five multiplies are necessary to compute this value (assuming that the value  $(v^2)^2$  is computed once and reused). More generally, raising a value to an integer power n can be done with  $O(\log n)$  multiplies.

An obvious (and indeed correct) approach to compute  $v^{20}$  would be to call std::pow(v, 20). Doing so is likely to be much less efficient than five multiplies; in general, std::pow() is implemented by computing a logarithm, multiplying by the exponent, and then exponentiating. Evaluating these transcendental functions generally requires tens or hundreds of machine instructions, taking much longer than a handful of multiplies. Some std::pow() implementations check for small integer exponents and handle them specially, and some compilers detect calls like std::pow(v, 2) and directly turn them into multiplies, but neither of these can be depended on.

C++ template functions offer a way to turn exponentiation like this into an efficient series of multiplies. Consider this use of templates:

```
template <int n> Float Pow(Float v) { return v * Pow<n-1>(v); }
template <> Float Pow<0>(Float v) { return 1; }
```

A call like Pow<5>(v) will be turned into v\*v\*v\*v\*v, which can be directly compiled into a series of multiplies. There's a catch, however: recall from Section 3.9 of the third edition that the IEEE floating-point standard prohibits the compiler from reassociating floating-point expressions. Thus, a call Pow<20>(v) will be compiled to nineteen multiply operations—likely more efficient than std::pow(), but not yet the logarithmic number that's possible assuming we don't mind reassociation.

The following template functions give us a logarithmic number of multiplication operations. The main Pow() function splits the exponent in half before recursively calling itself; template specializations handle the base cases.

```
⟨General Utility Functions⟩ ≡
template <int n>
static Float Pow(Float v) {
   static_assert(n > 0, "Power can't be negative");
   Float n2 = Pow<n / 2>(v);
   return n2 * n2 * Pow<n & 1>(v);
}
template <> Float Pow<1>(Float v) { return v; }
template <> Float Pow<0>(Float v) { return 1; }
```

In benchmarks on a 2016-era laptop, this implementation was 4.6x faster than calling std::pow for its uses later in this code.

Though our implementation of Pow() is straightforward, it's always a good idea to have a unit test. pbrt uses the Google Test framework for unit tests; see the pbrt-v3 User's Guide for more information. Here, we test integer powers up to 29. The value 2 for v is chosen carefully; recall from the discussion on p. 214 of the

third edition that with IEEE floating-point arithmetic, multiplication by a factor of two gives an exact result as long as there's no underflow or overflow. Thus, we can reasonably expect exact equality with the integer power of two reference values computed for the test.

```
{Hair Tests} ≡
TEST(Hair, Pow) {
    EXPECT_EQ(Pow<0>(2.f), 1 << 0);
    EXPECT_EQ(Pow<1>(2.f), 1 << 1);
    EXPECT_EQ(Pow<2>(2.f), 1 << 2);
        ⟨Test remainder of pow template powers to 29⟩
}</pre>
```

In the following, we'll need to compute the arcsine of various values. These values may be slightly outside the legal range [-1, 1] due to floating-point round-off error; the SafeASin() utility function handles clamping to this range, which makes calling code a bit cleaner. A runtime assertion makes sure that the value provided isn't too far out of the valid range.

```
⟨General Utility Functions⟩ ≡
inline Float SafeASin(Float x) {
   CHECK(x >= -1.0001 && x <= 1.0001);
   return std::asin(Clamp(x, -1, 1));
}</pre>
```

Similarly, we need to compute the square root of values that may be slightly negative due to round-off error; again, the clamp to the valid range is nice to have in a single place.

```
⟨General Utility Functions⟩ ≡
inline Float SafeSqrt(Float x) {
   CHECK_GE(x, -1e-4);
   return std::sqrt(std::max(Float(0), x));
}
```

## **1.3 SCATTERING FROM HAIR**

Geometric setting and utility functions in hand, we will now turn to discuss the general scattering behaviors that give hair its distinctive appearance and some of the assumptions that we'll make in the following.

Hair and fur have three main components:

- Cuticle: the outer layer, which forms the boundary with air. The cuticle's surface is a nested series of scales at a slight angle to the hair surface.
- Cortex: the next layer inside the cuticle. The cortex generally accounts for around 90% of hair's volume but less for fur. It is typically colored with pigments that mostly absorb light.
- Medulla: the center core at the middle of the cortex. It is larger and more significant in thicker hair and fur. The medulla is also pigmented. Scattering

**Figure 1.6:** The surface of hair is formed by scales that deviate a by a small angle 
$$\alpha$$
 from the idea cylinder. ( $\alpha$  is generally around  $2 - 4^{\circ}$ ; the angle shown here is larger for illustrative purposes.)

from the medulla is much more significant than scattering from the medium in the cortex.

For the following model, we'll make a few assumptions. (Approaches for relaxing some of them are discussed in the exercises at the end of this document.) First, we assume that the cuticle can be modeled as a rough dielectric cylinder with scales that are all angled at the same angle  $\alpha$  (effectively giving a nested series of cones.) (Figure 1.6.) We also treat the hair interior as a homogeneous medium that only absorbs light—scattering inside the hair is not modeled directly.

We will also make the assumption that scattering can be modeled accurately by a BSDF—we model light as entering and exiting the hair at the same place. (A BSSRDF could certainly be used instead; it's unclear how important subsurface light transport is in practice.) Note that this assumption does require that the hair's diameter be fairly small with respect to how quickly illumination changes over the surface; this assumption is generally fine in practice.

Incident light arriving at a hair may be scattered one more more times before leaving the hair; Figure 1.7 shows a few of the possible cases. We use p to denote the number of path segments it follows inside the hair before being scattered back out to air. We will sometimes refer to terms with a shorthand that describes the corresponding scattering events at the boundary: p = 0 corresponds to R, for reflection, p = 1 is TT, for two transmissions p = 2 is TRT, p = 3 is TRRT, and so forth.

In the following, we will find it useful to consider these scattering modes separately and so will write the hair BSDF as a sum over terms p:

$$f(\omega_{\rm o}, \omega_{\rm i}) = \sum_{p=0}^{\infty} f_p(\omega_{\rm o}, \omega_{\rm i}).$$
(1.1)

To make the scattering model implementation and importance sampling easier, many hair scattering models factor f into terms where one depends only on the angles  $\theta$  and another on  $\phi$ , the difference between  $\phi_0$  and  $\phi_i$ . This *semi-separable* model is given by:

$$f_p(\omega_{\rm o}, \omega_{\rm i}) = \frac{M_p(\theta_{\rm o}, \theta_{\rm i}) A_p(\omega_{\rm o}) N_p(\phi)}{|\cos \theta_{\rm i}|},$$
(1.2)

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where we have a longitudinal scattering function  $M_p$ , an attenuation function  $A_p$ , and an azimuthal scattering function  $N_p$ <sup>2</sup>. The division by  $|\cos \theta_i|$  cancels out the corresponding factor in the reflection equation.

In the following implementation, we will evaluate the first few terms of the sum in Equation (1.1) and then represent all higher-order terms with a single one. The pMax constant controls how many are evaluated before the switch-over.

 $\langle HairBSDF \ Constants \rangle \equiv$ static const int pMax = 3;

The model implemented in the HairBSDF is parameterized by six values:

- h: the [-1, 1] offset along the curve width where the ray intersected the oriented ribbon (h was defined in Section 1.1).
- eta: the index of refraction of the interior of the hair. (Typically, 1.55).
- sigma\_a: the absorption coefficient of the hair interior, where distance is measured with respect to the hair cylinder's diameter.
- •beta\_m: the longitudinal roughness of the hair, mapped to the range [0, 1].
- beta\_n: the azimuthal roughness, also mapped to [0, 1].
- alpha: the angle that the small scales on the surface of hair are offset from the base cylinder, expressed in degrees. (Typically, 2).

<sup>2</sup> Other authors generally include  $A_p$  in the  $N_p$  term, though we find it more clear to keep them separate for the following exposition. Here we also use f for the BSDF, which most hair scattering papers denote by S.

```
const Float h, gammaO, eta;
const Spectrum sigma_a;
const Float beta_m, beta_n, alpha;
```

We'll proceed to the method that evaluates the BSDF, leaving implementation of the code fragments in the constructor for later, closer to where the values they compute are used.

```
\langle HairBSDF Method Definitions \rangle \equiv
```

```
Spectrum HairBSDF::f(const Vector3f &wo, const Vector3f &wi) const {

\langle Compute \ hair \ coordinate \ system \ terms \ related \ to \ wo \rangle

\langle Compute \ hair \ coordinate \ system \ terms \ related \ to \ wi \rangle

\langle Compute \ cos \ \theta_t \ for \ refracted \ ray \rangle

\langle Compute \ \gamma_t \ for \ refracted \ ray \rangle

\langle Compute \ the \ transmittance \ T \ of \ a \ single \ path \ through \ the \ cylinder \rangle

\langle Evaluate \ hair \ BSDF \rangle

}
```

There are a few quantities related to the directions  $\omega_{\rm o}$  and  $\omega_{\rm i}$  that are needed for evaluating the hair scattering model—specifically, the sine and cosine of the angle  $\theta$  that each direction makes with the plane perpendicular to the curve, and the angle  $\phi$  in the azimuthal coordinate system.

As explained in Section 1.1,  $\sin \theta_0$  is given by the *x* component of  $\omega_0$  in the BSDF coordinate system. Given  $\sin \theta_0$ , because  $\theta_0 \in [-\pi/2, \pi/2]$ , we know that  $\cos \theta_0$  must be positive, and so we can compute  $\cos \theta_0$  using the identity  $\sin^2 \theta + \cos^2 \theta = 1$ . The angle  $\phi_0$  in the perpendicular plane can be computed with std::atan.

```
(Compute hair coordinate system terms related to wo) =
Float sinTheta0 = wo.x;
Float cosTheta0 = SafeSqrt(1 - Sqr(sinTheta0));
Float phi0 = std::atan2(wo.z, wo.y);
```

Equivalent code, not included here, computes these values for wi.

### **1.3.1 LONGITUDINAL SCATTERING**

Onward to  $M_p$ , the function that defines the component of scattering related to the angles  $\theta$ —longitudinal scattering. Longitudinal scattering is responsible for the specular lobe along the length of hair and the longitudinal roughness  $\beta_m$ controls the size of this highlight. Figure 1.8 shows a hair model rendered with three different longitudinal scattering roughnesses.

The model implemented here was developed from d'Eon et al. (2011). The mathematical details of the derivation are complex, so we won't include them here. The goals (which they achieved) were to derive a scattering function that is normalized (ensuring both energy conservation and no energy loss) and can be sampled directly. Although the model isn't derived based on a physical model of how hair scatters light, it matches measured data well and has parametric control of roughness v.

Their model is:

$$M_p(\theta_{\rm o}, \theta_{\rm i}) = \frac{1}{2v \sinh(1/v)} e^{-\frac{\sin \theta_{\rm i} \sin \theta_{\rm o}}{v}} I_0\left(\frac{\cos \theta_{\rm o} \cos \theta_{\rm i}}{v}\right), \qquad (1.3)$$

where  $I_0$  is the modified Bessel function of the first kind and v is the roughness variance. (Note that this is a different usage of v than earlier in this document when it was used for the parametric coordinate along the width of a curve.) Figure 1.9 shows plots of  $M_p$ .

It turns out that this model isn't numerically stable for low roughness variance values, so d'Eon (2013) derived a different approach for that case that operates on the log of  $I_0$  before taking an exponent at the end. The v <= .1 test in the implementation below selects between the two formulations.

10() and Log10() compute the values of the modified Bessel function of the first kind its logarithm, respectively. We won't include their implementations here, which are based on numerical approximations to those transcendental functions.

```
⟨Hair Local Declarations⟩ ≡
inline Float IO(Float x), LogIO(Float x);
```

One challenge with this model is choosing a roughness v to achieve a desired look. Here we have implemented a perceptually uniform mapping from roughness  $\beta_m \in [0, 1]$  to v where a roughness of 0 is nearly perfectly smooth and 1 is



Figure 1.8: The Effect of Varying the Longitudinal Roughness  $\beta_m$ . Hair model illuminated by a skylight environment map rendered with varying longitudinal roughness. (a) With a very low roughness,  $\beta_m = 0.1$ , the hair appears too shiny—almost metallic. (b) With  $\beta_m = 0.25$ , the highlight is similar to typical human hair. (c) At high roughness,  $\beta_m = 0.6$ , the hair is unrealistically flat and diffuse. (Hair geometry courtesy Cem Yuksel.)



extremely rough. Different roughness values are used for different values of p. For p = 1, roughness is reduced by an empirical factor that models the focusing of light due to refraction through the circular boundary of the hair. It is then increased for p = 2 and subsequent terms, which models the effect of light spreading out after multiple reflections at the rough cylinder boundary in interior of the hair. (See the "Further Reading" section for more on this variation.)

```
 \langle Compute \ longitudinal \ variance \ from \ \beta_m \rangle \equiv \\ v[0] = Sqr(0.726f * beta_m + 0.812f * Sqr(beta_m) + \\ 3.7f * Pow<20>(beta_m)); \\ v[1] = .25 * v[0]; \\ v[2] = 4 * v[0]; \\ for \ (int \ p = 3; \ p <= pMax; ++p) \\ v[p] = v[2]; \\ \langle HairBSDF \ Private \ Data \rangle \equiv \\ Float \ v[pMax + 1]; \end{cases}
```

### **1.3.2 ABSORPTION IN FIBERS**

The  $A_p$  term describes how much of the incident light is affected by each of the scattering modes p. It incorporates two effects: Fresnel reflection and transmission at the hair–air boundary and absorption of light that passes through the hair (for p > 0). This absorption is what gives hair and fur its color. Figure 1.10 has rendered images of hair with varying absorption coefficients, showing the effect that absorption has. The  $A_p$  function that we will implement here models all reflection and transmission at the hair boundary as perfectly specular—a very different assumption that  $M_p$  (and  $N_p$  to come), which model glossy reflection and transmission. This assumption simplifies the implementation and gives reasonable results in practice (presumably in that the specular paths are in a sense averages over all of the possibly glossy paths.)

We'll start by finding the transmittance of a single transmitted segment through the hair. To do so, we need to find the distance the ray travels until it exits the



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cylinder; the easiest way to do this is to compute the distances in the longitudinal and azimuthal projections separately.

To compute these distances, we need the transmitted angles of the ray  $\omega_0$ , in the longitudinal and azimuthal planes, which we'll denote by  $\theta_t$  and  $\gamma_t$ , respectively. Application of Snell's law using the hair's index of refraction  $\eta$  allows us to compute  $\sin \theta_t$  and  $\cos \theta_t$ .

 $\langle Compute \cos \theta_t \text{ for refracted ray} \rangle \equiv$ Float sinThetaT = sinThetaO / eta; Float cosThetaT = SafeSqrt(1 - Sqr(sinThetaT));

For  $\gamma_t$ , although we could compute the transmitted direction  $\omega_t$  from  $\omega_o$  and then project  $\omega_t$  into the normal plane, it's possible to compute  $\gamma_t$  directly using a *modified index of refraction* that accounts for the effect of the longitudinal angle on the refracted direction in the normal plane. The modified index of refraction is given by

$$\eta' = \frac{\sqrt{\eta^2 - \sin^2 \theta_{\rm o}}}{\cos \theta_{\rm o}}$$

Given  $\eta'$ , we can compute the refracted direction  $\gamma_t$  directly in the normal plane.<sup>3</sup> Since  $h = \sin \gamma_0$ , we can apply Snell's law (p. 546) to compute  $\gamma_t$ .

```
 \begin{array}{l} \langle Compute \; \gamma_t \; for \; refracted \; ray \rangle \equiv \\ & \mbox{Float etap = std::sqrt(eta * eta - Sqr(sinTheta0)) / cosTheta0;} \\ & \mbox{Float sinGammaT = h / etap;} \\ & \mbox{Float cosGammaT = SafeSqrt(1 - Sqr(sinGammaT));} \\ & \mbox{Float gammaT = SafeASin(sinGammaT);} \end{array}
```

If we consider the azimuthal projection of the transmitted ray in the normal plane, we can see that the segment makes the same angle  $\gamma_t$  with the circle normal at both of its endpoints (Figure 1.11). If we denote the total length of the segment by  $l_a$ , then basic trigonometry tells us that  $l_a/2 = \cos \gamma_t$ , assuming a unit radius circle.

Now considering the longitudinal projection, we can see that the distance that a transmitted ray travels before exiting is scaled by a factor of  $1/\cos\theta_t$  as it passes through the cylinder (Figure 1.12). Putting these together, the total segment length in terms of the hair diameter is

$$l = \frac{2\cos\gamma_{\rm t}}{\cos\theta_{\rm t}}$$

Recall that for the HairBSDF we defined  $\sigma_a$  to be measured with respect to the hair diameter (so that adjusting the hair geometry's width doesn't completely change its color). Therefore, we do not consider the hair cylinder diameter when

<sup>3</sup> This is due to the Bravais properties of cylindrical scattering. See Appendix B of Marschner et al. (2003) for a nice derivation and further explanation.





we apply Beer's law, and transmittance is given by

$$T = e^{-\sigma_{a}l}.$$
 (1.4)

Given a single segment's transmittance, we can now describe the function that evaluates the full  $A_p$  function. Ap() returns an array with the values of  $A_p$  up to  $p_{\max}$  and a final value that accounts for the sums of attenuations for all of the higher-order scattering terms.

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For the  $A_0$  term, corresponding to light that reflects at the cuticle, the Fresnel reflectance at the air-hair boundary gives the fraction of light that is reflected. We can find the cosine of the angle between the surface normal and the direction vector with angles  $\theta_0$  and  $\gamma_0$  in the hair coordinate system by  $\cos \theta_0 \cos \gamma_0$ .

```
 \langle Compute \ p = 0 \ attenuation \ at initial \ cylinder \ intersection \rangle \equiv \\ \texttt{Float } \cos\texttt{Gamma0} = \texttt{SafeSqrt(1 - h * h);} \\ \texttt{Float } \cosh\texttt{Theta} = \texttt{cosTheta0} * \texttt{cosGamma0;} \\ \texttt{Float } f = \texttt{FrDielectric(cosTheta, 1.f, eta);} \\ \texttt{ap[0]} = \texttt{f;} \end{aligned}
```

For the TT term, p = 1, we have two 1 - f terms, accounting for transmission into and out of the cuticle boundary, and a single T term for one transmission path through the hair. For all of the p > 0 terms, which include transmission, we can neglect the scaling of radiance based on the different indices of refraction of the exterior and interior media (recall the discussion of this effect on p. 527): because the viewer and light source are both assumed to be outside the hair, all of those factors cancel out.

⟨Compute p = 1 attenuation term⟩ ≡
ap[1] = Sqr(1 - f) \* T;

The p = 2 term has one more reflection event, reflecting light back into the hair, and then a second transmission term. Since we assume perfect specular reflection at the cuticle boundary, both segments inside the hair make the same angle  $\gamma_t$ with the circle's normal (Figure 1.13). From this, we can see that both segments must have the same length (and so forth for subsequent segments.) In general, for p > 0,

$$A_p = (1 - f)^2 T^p f^{p-1}.$$

{Compute attenuation terms up to p = pMax} =
for (int p = 2; p < pMax; ++p)
ap[p] = ap[p - 1] \* T \* f;</pre>

After pMax, a final term accounts for all further orders of scattering. We'd like to compute the sum of the infinite series of remaining terms, which fortunately can be found in closed form, since both T < 1 and f < 1:



$$\sum_{p=p_{\max}}^{\infty} (1-f)^2 T^p f^{p-1} = \frac{(1-f)^2 T^{p_{\max}} f^{p_{\max}-1}}{1-Tf}$$

(Compute attenuation term accounting for remaining orders of scattering) =
ap[pMax] = ap[pMax - 1] \* f \* T / (Spectrum(1.f) - T \* f);

## **1.3.3 AZIMUTHAL SCATTERING**

Finally, we will model the component of scattering dependent on the angle  $\phi$ . We will do this work entirely in the normal plane. The azimuthal scattering model is based on first computing a new azimuthal direction assuming perfect specular reflection and transmission and then defining a distribution of directions around this central direction, where increasing roughness gives a wider distribution. Therefore, we will first consider how an incident ray is deflected by specular reflection and transmission in the normal plane; Figure 1.14 illustrates the cases for the first two values of p.

Following the reasoning from Figure 1.14, we can derive the function  $\Phi$ , which gives the net change in azimuthal direction:

$$\Phi(p,h) = 2p\gamma_{\rm t} - 2\gamma_{\rm o} + p\pi.$$

(Recall that  $\gamma_0$  and  $\gamma_t$  are derived from h.) Figure 1.15 shows a plot of this function for p = 1.



angle  $\gamma_{\rm o}$ , with the surface normal. The net change in angle is thus  $-2\gamma_{\rm o}$ . For p=1, the ray is deflected from  $\gamma_{\rm o}$  to  $\gamma_{\rm t}$  when it enters the cylinder and then correspondingly on the way out. We can also see that when the ray is transmitted again out of the circle, it again makes an angle  $\gamma_{\rm o}$  with the surface normal there. Adding up the angles, the net deflection is  $2\gamma_{\rm t}-2\gamma_{\rm o}+\pi.$ 

```
{Hair Local Functions} =
inline Float Phi(int p, Float gamma0, Float gammaT) {
    return 2 * p * gammaT - 2 * gamma0 + p * Pi;
}
```

Now that we know how to compute new angles in the normal plane after specular transmission and reflection, we need a way to represent surface roughness, so that a range of directions centered around the specular direction can contribute to scattering. The *logistic distribution* provides a good option: it is a generally useful one for rendering, since it has a similar shape to the Gaussian (which of course comes up often in rendering), while also being normalized and integrable in closed-form (unlike the Gaussian).

The logistic distribution takes a scale factor s, which controls its width:

$$l(x,s) = \frac{e^{-x/s}}{s(1 + e^{-x/s})^2}$$

The implementation is straightforward, though it is worth taking the absolute value of x to avoid numerical instability for when the ratio x/s is relatively



large. (The function is symmetric around the origin, so this is mathematically equivalent.)

```
(Hair Local Functions) =
inline Float Logistic(Float x, Float s) {
    x = std::abs(x);
    return std::exp(-x / s) / (s * Sqr(1 + std::exp(-x / s)));
}
```

Because the logistic distribution is normalized, it is its own PDF. Its integral is reasonably straightforward:

$$\int l(x,s) \, \mathrm{d}x = \frac{1}{1 + \mathrm{e}^{-x/s}},\tag{1.5}$$

and the function that implements its CDF follows directly.

```
(Hair Local Functions) =
  inline Float LogisticCDF(Float x, Float s) {
    return 1 / (1 + std::exp(-x / s));
}
```

In the following, we'll find it useful to define a normalized logistic function over a range [a, b]; we'll call this the *trimmed logistic*,  $l_t$ . (In practice, we'll always use the range  $[-\pi, \pi]$ , but will derive the next few functions for arbitrary ranges for flexibility.)

$$l_{t}(x, s, [a, b]) = \frac{l(x, s)}{\int_{a}^{b} l(x', s) \,\mathrm{d}x'}$$

The implementation follows directly using Equation (1.5). Figure 1.16 shows plots of the trimmed logistic distribution for a few values of s.





```
\langle Hair \ Local \ Functions \rangle \equiv
```

```
inline Float TrimmedLogistic(Float x, Float s, Float a, Float b) {
    return Logistic(x, s) / (LogisticCDF(b, s) - LogisticCDF(a, s));
}
```

Now we have the pieces to be able to implement the azimuthal scattering distribution. The Np() function computes the  $N_p$  term, computing the angular difference between  $\phi$  and  $\Phi(p, h)$  and evaluating the azimuthal distribution with that angle.

The difference between  $\phi$  and  $\Phi(p, h)$  may be outside the range we've defined the logistic over,  $[-\pi, \pi]$ , so we rotate around the circle as needed to get the value to the right range. Because dphi never gets too far out of range for the small p used here, we just use the simple approach of adding or subtracting  $2\pi$  as needed.

 $\langle Remap \text{ dphi } to [-\pi, \pi] \rangle \equiv$ while (dphi > Pi) dphi -= 2 \* Pi; while (dphi < -Pi) dphi += 2 \* Pi;

As with the longitudinal roughness, it's helpful to have a roughly perceptually linear mapping from azimuthal roughness  $\beta_n \in [0, 1]$  to the logistic scale factor s.

 $\langle Compute \ azimuthal \ logistic \ scale \ factor \ from \ \beta_n \rangle \equiv \\ s = SqrtPiOver8 * (0.265f * beta_n + 1.194f * Sqr(beta_n) + \\ 5.372f * Pow<22>(beta_n));$ 



The mapping uses the constant  $\sqrt{\pi/8}$ .

 $\langle HairBSDF \ Constants \rangle \equiv \\ \texttt{static const Float SqrtPiOver8 = 0.626657069f;}$ 

 $\langle HairBSDF Private Data \rangle \equiv$ Float s;

Figure 1.17 shows polar plots of azimuthal scattering for the TT term, p = 1, with a fairly low roughness. The scattering distributions for the two different points on the curve's width are quite different. Because we expect the hair width to be roughly pixel-sized, many rays per pixel are needed to resolve this variation well.

Figure 1.18 shows renderings of the hair model from Figure 1.1 with a fixed longitudinal roughness and varying azimuthal roughness. We can see that higher azimuthal roughness causes the hair to be lighter in color; this is because more light is able to exit the hair after multiple scattering when the distribution is broader.

## **1.3.4 SCATTERING MODEL IMPLEMENTATION**

We now have almost all of the pieces we need to be able to evaluate the model. The last detail is to account for the effect of scales on the hair surface (recall Figure 1.6). Suitable adjustments to  $\theta_i$  work well to model this characteristic of hair.

For the R term, adding the angle  $2\alpha$  to  $\theta_i$  can model the effect of evaluating the hair scattering model with respect to the surface normal of a scale. We can then go ahead and evaluate  $M_0$  with this modification to  $\theta_i$ . For TT, we have to account for two transmission events through scales. Rotating by  $\alpha$  in the opposite direction approximately compensates. (Because the refraction angle is





non-linear with respect to changes in normal orientation, there is some error in this approximation, though the error is low for the typical case of small values of  $\alpha$ .) TRT has a reflection term inside the hair; a rotation by  $-4\alpha$  compensates for the overall effect.

The effects of these shifts are that the primary reflection lobe R is offset to be above the perfect specular direction and the secondary TRT lobe is shifted below it. Together, these lead to two distinct specular highlights of different colors, since R isn't affected by the hair's color, while TRT picks up the hair color due to absorption. This effect can be seen in human hair. Figure 1.19 shows the visual result of accounting for the tilted scales.

To support computing these offsets, in the HairBSDF constructor we precompute  $\sin 2^k \alpha$  and  $\cos 2^k \alpha$  for k = 0, 1, 2. These values can be computed particularly efficiently using trigonometric double angle identities:  $\cos 2\theta = \cos^2 \theta - \sin^2 \theta$  and  $\sin 2\theta = 2 \cos \theta \sin \theta$ .

```
(Compute α terms for hair scales) =
sin2kAlpha[0] = std::sin(alpha);
cos2kAlpha[0] = SafeSqrt(1 - Sqr(sin2kAlpha[0]));
for (int i = 1; i < 3; ++i) {
    sin2kAlpha[i] = 2 * cos2kAlpha[i - 1] * sin2kAlpha[i - 1];
    cos2kAlpha[i] = Sqr(cos2kAlpha[i - 1]) - Sqr(sin2kAlpha[i - 1]);
}</pre>
```

⟨HairBSDF Private Data⟩ ≡
Float sin2kAlpha[3], cos2kAlpha[3];

#### SECTION 1.3 SCATTERING FROM HAIR

Evaluating the model is now mostly just a matter of calling functions that have already been defined and summing the individual terms  $f_p$ .

As discussed earlier,  $\theta_i$  is rotated to model the effect of hair scales. Fortunately we only need the sine and cosine of the angle  $\theta_i$  to evaluate  $M_p$ . We can therefore use the trigonometric identities

 $\sin \theta \pm \alpha = \sin \theta \cos \alpha \pm \cos \theta \sin \alpha$  $\cos \theta \pm \alpha = \cos \theta \cos \alpha \mp \sin \theta \sin \alpha$ 

to efficiently compute the rotated angle, without needing to evaluate any additional trigonometric functions.

Here we only include the case for p = 0, where  $\theta_i$  is rotated by  $2\alpha$ . The remaining cases follow the same structure. (For p = 1, the rotation is by  $-\alpha$  and for p = 2,  $-4\alpha$ .)

 $\langle Compute \sin \theta_i \text{ and } \cos \theta_i \text{ terms accounting for scales} \rangle \equiv$ 

Float sinThetaIp, cosThetaIp; if (p == 0) { sinThetaIp = sinThetaI \* cos2kAlpha[1] + cosThetaI \* sin2kAlpha[1]; cosThetaIp = cosThetaI \* cos2kAlpha[1] - sinThetaI \* sin2kAlpha[1]; } {Handle remainder of p values for hair scale tilt} {Handle out-of-range cos θ<sub>i</sub> from scale adjustment}

When  $\omega_i$  is nearly parallel with the hair, the scale adjustment may give a slightly negative value for  $\cos \theta_i$ —effectively, in this case, it represents a  $\theta_i$  that is slightly greater than  $\pi/2$ , the maximum expected value of  $\theta$  in the hair coordinate system. This angle is equivalent to  $\pi - \theta_i$ , and  $\cos(\pi - \theta_i) = |\cos \theta_i|$ , so we can easily handle that here.

 $\langle Handle \ out\text{-}of\text{-}range \ \cos \theta_i \ from \ scale \ adjustment \rangle \equiv \\ \texttt{cosThetaIp} = \texttt{std::abs(cosThetaIp);}$ 

A final term accounts for all higher-order scattering inside the hair. We just use a uniform distribution  $N(\phi) = 1/(2\pi)$  for the azimuthal distribution; this is a reasonable choice, as the varied direction offsets from  $\Phi(p, h)$  for  $p \ge p_{\text{max}}$  generally have wide variation and the final  $A_p$  term generally represents less than 15% of the overall scattering, so little error is introduced in the final result.

(Compute contribution of remaining terms after pMax) =
fsum += Mp(cosThetaI, cosThetaO, sinThetaI, sinThetaO, v[pMax]) \*
ap[pMax] / (2.f \* Pi);

## 1.3.5 THE "WHITE FURNACE" TEST

We have, we hope, implemented a hair scattering model wherein if hair doesn't absorb any of the light passing through it (i.e.,  $\sigma_a = 0$ ), then all of the incident light should be reflected. If such a hair is illuminated with uniform incident radiance, the reflected radiance should be exactly the same as the incident radiance. The *white furnace test* checks this, making sure that reflected radiance is one given unit incident radiance. Our implementation tests a variety of azimuthal and longitudinal roughnesses.

For each roughness, we compute a Monte Carlo estimate of the sphericaldirectional reflectance,

$$\int_{\mathbb{S}^2} f(\omega_{\rm o}, \omega_{\rm i}) \left| \cos \theta_{\rm i} \right| \, \mathrm{d}\omega_{\rm i}.$$

Each sample is evaluated by first sampling a random offset along the hair h and then computing the fraction of reflected radiance for a random incident direction.

 $\langle Estimate \ reflected \ uniform \ incident \ radiance \ from \ hair \rangle \equiv$ 

#### SECTION 1.4 IMPORTANCE SAMPLING

The WhiteFurnaceSampled test, not included here, uses the Sample\_f() method (which will be defined shortly) to sample incident directions rather than a uniform distribution, dividing BSDF values by the PDF to compute the estimate of reflectance. It uses a tighter tolerance  $(1 \pm .01)$  than the first white furnace test, since importance sampling gives much faster convergence than uniform sampling over the sphere. (Note, however, that it's useful to have both variants of this test: if there was a bug in the importance sampling code and the second white furnace test failed, we wouldn't know whether the bug was in the sampling code or how the hair BSDF computed reflection. This way, if one white furnace test fails or both fail, we can have a better idea of whether the underlying problem is in the evaluation of the model or in the code that samples it.)

## **1.4 IMPORTANCE SAMPLING**

Being able to generate sampled directions and compute the PDF for sampling a given direction according to a distribution that is similar to the overall BSDF is critical for efficient rendering, especially at low roughnesses, where the hair BSDF varies rapidly as a function of direction. In the approach implemented here, samples are generated with a two step process: first we choose a p term to sample according to a probability based on each term's  $A_p$  function value, which gives its contribution to the overall scattering function. Then, we find a direction by sampling the corresponding  $M_p$  and  $N_p$  terms. Fortunately, both the  $M_p$  and  $N_p$  terms of the hair BSDF can be sampled perfectly, leaving us with a sampling scheme that exactly matches the PDF of the full BSDF.

#### **1.4.1 COMPUTING ADDITIONAL SAMPLE VALUES**

In the following, we'll need a total of four random samples to sample the direction wi. This presents a challenge in that only two sample values are passed into pbrt's BxDF::Sample\_f() interface. One option would be to modify the interfaces to provide more sample values, though these would be unused by all of the other BxDF implementations; rendering efficiency would suffer from the time to generate these unused samples, and rendering quality would also likely suffer, as many Samplers generate better samples in the lower dimensions than in higher dimensions; these extra samples would usually be wasted.

Therefore, here we'll implement an approach that lets us extract two separate samples from each provided sample. The DemuxFloat() function, to be defined shortly, decomposes a sample  $\xi \in [0, 1)$  into a pair of samples, while also making some effort to preserve stratification in the returned sample value.

To understand the operation of DemuxFloat(), first recall the discussion of the Morton curve in Section 4.3.3. The Morton curve is a space-filling 1D curve that maps real numbers in [0, 1] to *n*-dimensional numbers  $[0, 1]^n$ . If we use a 1D sample value  $\xi$  as an offset into a Morton curve, we can use each of the dimensions' values as independent samples.

An important advantage of using a Morton curve for this task is that it preserves stratification in the sample values it returns. To see why this is so, consider a collection of four stratified 1D sample values (i.e., one in [0, 1/4), one in [1/4, 1/2), and so forth.) The first sample value will be mapped mapped to a point in  $[0, 1/2) \times [0, 1/2)$  by the Morton curve, since the first quarter of the 2D Morton curve traces out that region of space in two dimensions. Similarly, the next stratified 1D sample will be in the range  $[1/2, 1) \times [0, 1/2)$ , and so on.

To help with this, the Compact1By1() function takes a 32-bit integer, deletes all of the bits with odd indices, and compacts the remaining bits. The comment lines in the implementation illustrate the effect of each operation: after the surviving bits are initially masked off, a series of masked shifts compacts them until they are in contiguous positions.<sup>4</sup>

```
\langle General \ Utility \ Functions \rangle \equiv
```

```
static uint32_t Compact1By1(uint32_t x) {
    // x = -f-e -d-c -b-a -9-8 -7-6 -5-4 -3-2 -1-0
    x &= 0x5555555;
    // x = --fe --dc --ba --98 --76 --54 --32 --10
    x = (x ^ (x >> 1)) & 0x33333333;
    // x = ---- fedc ---- ba98 ---- 7654 ---- 3210
    x = (x ^ (x >> 2)) & 0x0f0f0f0f;
    // x = ---- fedc ba98 ---- 7654 3210
    x = (x ^ (x >> 4)) & 0x00ff00ff;
    // x = ---- fedc ba98 7654 3210
    x = (x ^ (x >> 8)) & 0x0000ffff;
    return x;
}
```

Here, we'll treat the sample  $\xi$  as a fixed-point value v computed by multiplying by  $2^{32}$ . The 2D Morton curve effectively takes alternating bits from this value, giving two values between 0 and  $2^{16}$ . In turn, these are mapped back to Floats.

```
(General Utility Functions) =
   static Point2f DemuxFloat(Float f) {
        uint64_t v = f * (1ull << 32);
        uint32_t bits[2] = {Compact1By1(v), Compact1By1(v >> 1)};
        return {bits[0] / Float(1 << 16), bits[1] / Float(1 << 16)};
}</pre>
```

## 1.4.2 A DISTRIBUTION FOR SAMPLING P

Next, we'll define the ComputeApPdf() method, which returns a discrete PDF with probabilities for sampling each term  $A_p$  according to its contribution relative to all of the  $A_p$  terms, given  $\theta_0$ .

<sup>4</sup> This code is thanks to Fabian Giesen, fgiesen.wordpress.com/2009/12/13/decoding-morton-codes/.

The method starts by computing the values of  $A_p$  for cosThetaO. We're able to reuse some previously-defined fragments to make this task easier.

```
 \begin{array}{l} \langle Compute \ array \ of \ A_p \ values \ for \ cosTheta0 \rangle \equiv \\ \mbox{Float sinTheta0 = SafeSqrt(1 - cosTheta0 * cosTheta0);} \\ \langle Compute \ cos \ \theta_t \ for \ refracted \ ray \rangle \\ \langle Compute \ \gamma_t \ for \ refracted \ ray \rangle \\ \langle Compute \ the \ transmittance \ T \ of \ a \ single \ path \ through \ the \ cylinder \rangle \\ \mbox{std::array<Spectrum, pMax + 1> ap = Ap(cosTheta0, eta, h, T);} \end{array}
```

Next, the spectral  $A_p$  values are converted to scalars using their luminance and these values are normalized to make a proper PDF.

```
 \begin{array}{l} \langle Compute \; A_p \; PDF \; from \; individual \; A_p \; terms \rangle \equiv \\ \texttt{std::array<Float, pMax + 1> apPdf;} \\ \texttt{Float sumY = std::accumulate(ap.begin(), ap.end(), Float(0),} \\ [](\texttt{Float s, const Spectrum &ap}) \; \{ \; \texttt{return s + ap.y(); } \}); \\ \texttt{for (int i = 0; i <= pMax; ++i)} \\ & \texttt{apPdf[i] = ap[i].y() / sumY;} \end{array}
```

## **1.4.3 SAMPLING INCIDENT DIRECTIONS**

With these preliminaries out of the way, we can now implement the Sample\_f() method.

⟨Derive four random samples from u2⟩ ≡
Point2f u[2] = { DemuxFloat(u2[0]), DemuxFloat(u2[1]) };

Given the PDF over  $A_p$  terms, we just loop over PDF values until we find the first value of p where the sum of preceding PDF values is greater than the sample value. Because we only need to generate one sample from the PDF's distribution, the

work to compute an explicit CDF array (for example, by using  $\tt Distribution1D$ ) isn't worthwhile.

```
{Determine which term p to sample for hair scattering} =
std::array<Float, pMax + 1> apPdf = ComputeApPdf(cosThetaO);
int p;
for (p = 0; p < pMax; ++p) {
    if (u[0][0] < apPdf[p]) break;
    u[0][0] -= apPdf[p];
}</pre>
```

Now that we've chosen a term, we can sample the corresponding  $M_p$  term given  $\theta_0$  to find  $\theta_i$ . The derivation of this sampling method is fairly involved, so we'll neither include the derivation nor the implementation here. This fragment,  $\langle Sample M_p \text{ to compute } \theta_i \rangle$ , consumes both of the sample values u[1][0] and u[1][1] and initializes variables sinThetaI and cosThetaI according to the sampled direction. After sampling a direction  $\theta_i$ , this fragment then applies the inverse of the rotation that will later be used to account for hair scales when the BSDF is evaluated.

Next we'll sample the azimuthal distribution  $N_p$ . For terms up to  $p_{\text{max}}$ , we take a sample from the logistic distribution centered around the exit direction given by  $\Phi(p, h)$ . For the last term, we sample from a uniform distribution.

```
\begin{split} &\langle Sample \; N_p \; to \; compute \; \Delta \phi \rangle \equiv \\ &\langle Compute \; \gamma_t \; for \; refracted \; ray \rangle \\ & \text{Float dphi;} \\ & \text{if } (p < p\text{Max}) \\ & \text{dphi} = \text{Phi}(p, \; \text{gamma0}, \; \text{gammaT}) \; + \\ & \quad \text{SampleTrimmedLogistic(u[0][1], s, -Pi, Pi);} \\ & \text{else} \\ & \quad \text{dphi} = 2 \; * \; \text{Pi} \; * \; u[0][1]; \end{split}
```

By inverting the CDF of the trimmed logistic, we can derive the recipe to sample from its distribution given a random variable  $\xi \in [0, 1)$ :

$$\begin{split} \xi &= \int_{a}^{x} l_{t}(x', s, [a, b]) \, \mathrm{d}x' \\ &= \frac{1}{\int_{a}^{b} l(x, s) \, \mathrm{d}x} \int_{a}^{x} l(x', s) \, \mathrm{d}x' \\ &= \frac{1}{\int_{a}^{b} l(x, s) \, \mathrm{d}x} \left(\frac{1}{1 + \mathrm{e}^{-x/s}} - \frac{1}{1 + \mathrm{e}^{-a/s}}\right) \end{split}$$

With a bit of algebra, we can solve for x:

$$x = -s \log \left( \frac{1}{\xi \int_a^b l(x,s) \,\mathrm{d}x + \frac{1}{1 + \mathrm{e}^{-a/s}}} - 1 \right).$$

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#### SECTION 1.4 IMPORTANCE SAMPLING

In practice, due to floating-point round-off, the implementation may compute an infinite value when  $\xi \to 1$ ; a clamp at the end ensures that the returned value is in the range [a, b].

Given  $\theta_i$  and  $\phi_i$ , we can compute the sampled direction wi. The math is similar to the SphericalDirection() function defined on p. 346, but with two important differences. First, because here  $\theta$  is measured with respect to the plane perpendicular to the cylinder rather than the cylinder axis, we need to compute  $\cos(\pi/2 - \theta) = \sin \theta$  for the coordinate with respect to the cylinder axis instead of  $\cos \theta$ . Second, because the hair shading coordinate system's  $(\theta, \phi)$  coordinates are oriented with respect to the +x axis, the order of dimensions passed to the Vector3f constructor is adjusted correspondingly, since the direction returned from Sample\_f() should be in the BSDF coordinate system.

Because we could sample directly from the  ${\cal M}_p$  and  ${\cal N}_p$  distributions, the overall PDF is

$$\sum_{p=0}^{p_{\max}} M_p(\theta_{\rm o}, \theta_{\rm i}) \tilde{A}_p(\omega_{\rm o}) N_p(\phi),$$

where  $\tilde{A}_p$  are the normalized luminance-weighted PDF terms. Note that  $\theta_i$  must be shifted to account for hair scales when evaluating the PDF; this is done in the same way (and with the same code fragment) as when the BSDF was evaluated.

 $\langle Compute \ PDF \ for \ sampled \ hair \ scattering \ direction \ wi \rangle \equiv$ 

The HairBSDF::Pdf() method performs the same computation was we just implemented for Sample\_f(). Therefore, the implementation isn't included here.

## **1.4.4 TESTING SAMPLING ROUTINES**

Because the sampling routine we have implemented exactly matches the PDF of the underlying BSDF, if we generate samples from the BSDF using Sample\_f(), then the ratio between the BSDF value and the PDF computed for this direction should be one (as long as there is no absorption in the hair). The SamplingWeights test checks this for a variety of roughnesses and random sample values.

Performing the test is mostly a matter of setting up enough context to call Sample\_f() and then verifying the ratio of the BSDF and the PDF.

```
\langle Check \; \texttt{HairBSDF::Sample_f()} \; sample \; weight \rangle \equiv
```

Note that we accept a small amount of error, accepting values that are close to one but not exactly equal to it, in order to allow for floating-point round-off error.

```
⟨Verify that hair BSDF sample weight is close to 1 for wi⟩ ≡
EXPECT_GT(f.y() * AbsCosTheta(wi) / pdf, 0.999);
EXPECT_LT(f.y() * AbsCosTheta(wi) / pdf, 1.001);
```

Another useful test is based on computing reflected radiance from a varying incident radiance function with the scattering equation, Equation (5.9) in the third edition. Given a sufficient number of samples, we should get the same result both if we use the custom importance sampling scheme we have implemented and if we use a uniform distribution of directions over the unit sphere. This case is tested with the SamplingConsistency test.

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The Li lambda function defines an incident radiance function with some (but not too much) variation as a function of direction. We keep this function fairly simple so that sampling the BSDF alone works well to compute reflected radiance and we can avoid the complexity of implementing multiple importance sampling in the test here.

```
⟨Declare variables for hair sampling test⟩ ≡
const int count = 64*1024;
Spectrum sigma_a = .25;
Vector3f wo = UniformSampleSphere({rng.UniformFloat(), rng.UniformFloat()});
auto Li = [](const Vector3f &w) -> Spectrum {
    return w.z * w.z;
};
Spectrum fImportance = 0, fUniform = 0;
```

For each sample in the Monte Carlo estimate, we choose a random point on the hair and use a pair of random numbers to sample an incident direction. We then use the regular Monte Carlo estimator to compute estimates using both sampling approaches.

 $\langle Compute \ estimates \ of \ scattered \ radiance \ for \ hair \ sampling \ test \rangle \equiv$ 

```
Float h = -1 + 2 * rng.UniformFloat();
HairBSDF hair(h, 1.55, sigma_a, beta_m, beta_n, 0.f);
Vector3f wi;
Float pdf;
Point2f u = {rng.UniformFloat(), rng.UniformFloat()};
Spectrum f = hair.Sample_f(wo, &wi, u, &pdf, nullptr);
if (pdf > 0) fImportance += f * Li(wi) * AbsCosTheta(wi) / (count * pdf);
wi = UniformSampleSphere(u);
fUniform += hair.f(wo, wi) * Li(wi) * AbsCosTheta(wi) /
(count * UniformSpherePdf());
```

In the end, the two estimates should be very close. Here we treat 5% relative error as good enough to pass the test; this is a fairly low bar, but it allows the test to run quickly—if the total number of samples count was higher, we could expect closer agreement, but we prefer to have a test that runs in a second or so.

```
\langle Verify consistency of estimated hair reflected radiance values \rangle =
Float err = std::abs(fImportance.y() - fUniform.y()) / fUniform.y();
EXPECT_LT(err, 0.05);
```

## **1.5 HAIR ABSORPTION COEFFICIENTS**

The color of hair is determined by how pigments in the cortex absorb light, which in turn is described by the normalized absorption coefficient where distance is measured in terms of the hair diameter. If a specific hair color is desired, there's a non-obvious relationship between the normalized absorption coefficient and the color of hair in a rendered image. Not only does changing the spectral values of the absorption coefficient have an unpredictable connection to the appearance of a single hair, but as we saw in Figure 1.2, multiple scattering between collections of many hairs has a significant effect each one's apparent color.<sup>5</sup> Therefore, here we provide implementations of two more intuitive ways to specify hair color.

The color of human hair is determined by the concentration of two pigments. The concentration of eumelanin is the primary factor that causes the difference between black, brown, and blonde hair. (Black hair has the most eumelanin and blonde hair has the least. White hair has none.) The second pigment, pheomelanin, causes hair to be orange or red. The HairBSDF class provides a convenience method that computes an absorption coefficient using the product of user-supplied pigment concentrations and absorption coefficients of the pigments computed by d'Eon et al. (2011), based on a model by Donner and Jensen (2006).

```
⟨HairBSDF Method Definitions⟩ ≡
Spectrum HairBSDF::SigmaAFromConcentration(Float ce, Float cp) {
Float sigma_a[3];
Float eumelaninSigmaA[3] = {0.419f, 0.697f, 1.37f};
Float pheomelaninSigmaA[3] = {0.187f, 0.4f, 1.05f};
for (int i = 0; i < 3; ++i)
            sigma_a[i] = (ce * eumelaninSigmaA[i] +
                cp * pheomelaninSigmaA[i]);
return Spectrum::FromRGB(sigma_a);
}</pre>
```

Eumelanin concentrations of roughly 8, 1.3, and 0.3 give reasonable representations of black, brown, and blonde hair, respectively.

It's also useful to specify the desired hair color directly. In order to make this possible, Chiang et al. (2016) created a cube of hair and rendered it with a variety of absorption coefficients and roughnesses, while it was illuminated with a uniform white dome. They then fit a function that mapped from the hair's azimuthal roughness and average color at the center of the front face

<sup>5</sup> These issues are both related to those that make it difficult to directly set scattering and absorption coefficients for subsurface scattering, as discussed along with the SubsurfaceFromDiffuse() function defined on p. 938.

#### SECTION 1.6 HAIR MATERIAL



of the cube to an absorption coefficient. (Unlike the azimuthal roughness, the longitudinal roughness doesn't meaningfully affect the hair's color.) This function is implemented in the SigmaAFromReflectance() method; see Figure 1.20 for examples.

## **1.6 HAIR MATERIAL**

Like most Materials in pbrt, the HairMaterial is pretty straightforward; it's mostly a matter of evaluating textures and creating a corresponding HairBSDF

object. It allows the attenuation coefficient for the hair interior to be specified with one of three ways; the corresponding parameters are:

- "spectrum sigma\_a": the absorption coefficient can be specified directly.
- "spectrum reflectance": a reflectance for SigmaAFromReflectance().
- "float eumelanin" and/or "float pheomelanin": eumelanin and pheomelanin concentrations.

Only one of these approaches can be used; an error is issued if more than one is provided. If none is specified, an eumelanin concentration of 1.3 is used, giving a brownish color.

A few additional parameter are supported

- "float beta\_m", "float beta\_n": longitudinal and azimuthal roughnesses. Both default to 0.3.
- "float eta": index of refraction of the hair interior. (1.55 by default).
- "float alpha": hair scale angle in degrees (2 by default).

## **1.7 A NOTE ON RECIPROCITY**

On p. 350 in the third edition of *Physically Based Rendering*, we noted that physically-based BRDFs are both reciprocal and energy conserving. (In particular, reciprocity means that swapping the two evaluation directions gives the same BRDF value:  $f(\omega_0, \omega_i) = f(\omega_i, \omega_0)$ .) BTDFs are in general not reciprocal, however; this topic (and methods to address it) is discussed further on p. 960 when bidirectional light transport algorithms are introduced.

The model we have implemented is, unfortunately, not reciprocal. One immediate issue is that the rotation for hair scales is applied only to  $\theta_i$ . However, there are more problems: first, all terms p > 0 that involve transmission are not reciprocal; the underlying issue is that the transmission terms use values based on  $\omega_t$ , which itself only depends on  $\omega_0$ . Thus, if  $\omega_0$  and  $\omega_i$  are interchanged, a completely different  $\omega_t$  is computed, which in turn leads to different  $\cos \theta_t$  and  $\gamma_t$  values, which in turn give different values from the  $A_p$  and  $N_p$  functions.

Many earlier hair scattering models have worked around the lack of reciprocity due to  $\cos \theta_t$  by computing an angle  $\theta_d = (\theta_o - \theta_i)/2$  and using that in place of  $\theta_o$  when computing  $\theta_t$  and related quantities. (Note that  $\theta_d$  is symmetric, since it's measuring angles with respect to the normal plane.) We didn't use  $\theta_d$  in our implementation for two reasons: first, when we're sampling the BSDF, only  $\omega_o$ is known and thus  $\theta_d$  can't be computed. In turn,  $\theta_o$  must be used to generate the PDF of  $A_p$  terms. Later, when  $\omega_i$  is known and the BSDF is evaluated, a different value of  $A_p$  would be computed than was used for sampling. This mismatch between function value and PDF can cause a variance spike.<sup>6</sup>

<sup>6</sup> Previous approaches have addressed this issue by clamping the maximum ratio of function value and PDF, though doing so loses energy.

Second, even if  $\theta_d$  is used, there is a second, subtle and more difficult, source of non-reciprocity. Recall that the position h along the curve width is computed based on finding the intersection of a ray with a ribbon oriented to face the ray. It is thus directly dependent on  $\omega_0$  and independent of  $\omega_i$ . In turn, because  $\gamma_0$ and  $\gamma_t$  depend on h, they depend on  $\omega_0$  alone and thus reciprocity is lost. The models developed by Marschner et al. (2003) and d'Eon et al. (2011) didn't have this problem since both computed a BCSDF by integrating h across the curve width. As noted by Chiang et al. (2016), this integration is computationally expensive, especially with low azimuthal roughnesses where the function varies quickly and many evaluations are needed for deterministic quadrature methods (recall Figure 1.17).

Another option would be to integrate stochastically, sampling a random h. (In turn, any given evaluation of f wouldn't be reciprocal, but the expected value over a sum of many evaluations would be.) This approach doesn't fit well with pbrt's current architecture, where no random sample values are made available to the BSDF::f() method. Changing this would require a fairly extensive set of code modifications and seems generally unappealing.

A final possibility would be to compute a  $h_i$  for  $\omega_i$ , based on projecting the intersection point up to the surface of the cylinder and then finding  $h_i$  for a ray with direction  $\omega_i$  that passes through the hit point. In turn, we could evaluate  $A_p$  and  $N_p$  twice for each ordering of directions and take the average. This is relatively straightforward (and  $h_i$  can be found working entirely in the azimuthal plane), but in turn may lead to noise spikes, as one of the evaluations may have a much larger value than the current implementations of the sampling and PDF methods expect.

Stuck for an elegant solution, we will leave this issue to an exercise and hope that future research on this topic addresses this issue. In practice, we haven't seen visual artifacts in rendered images from the lack of reciprocity.

## **1.8 FURTHER READING**

Kajiya and Kay (1989) were the first to develop a hair scattering model for rendering. Their model combined a diffuse term with a Phong lobe for an empirical model of specular highlights.

Marschner et al. (2003) investigated the processes underlying scattering from hair and performed a variety of measurements of scattering from actual hair. They introduced the longitudinal/azimuthal decomposition and the use of the modified index of refraction to hair rendering. They then developed a scattering model where the longitudinal component was derived by first considering perfectly specular paths and then allowing roughness by centering a Gaussian around them, and their azimuthal model assumed perfectly specular reflections. They showed that this model agreed reasonably well with their measurements.

Zinke and Weber (2007) formalized different ways of modeling scattering from hair and clarified the assumptions underlying each of them. Starting with the bidirectional fiber scattering distribution function (BFSDF), which describes reflected differential radiance at a point on a hair as a fraction of incident differential power at another, they showed how assuming homogeneous scattering properties and a far away viewer and illumination made it possible to simplify the eight-dimensional BFSDF to a four-dimensional *bidirectional curve scattering distribution* (BCSDF).

d'Eon et al. (2011) made a number of improvements to Marschner et al.'s model. They showed that their  $M_p$  term wasn't actually energy conserving and derived a new one that was; this is the model from Equation (1.3) that our implementation uses. (See also d'Eon (2013) for a more numerically stable formulation of  $M_p$  for low roughness.) They also introduced a Gaussian to the azimuthal term, allowing for varying azimuthal roughness. A 1D quadrature method was used to integrate the model across the width of the hair h.

Being able to generate samples from a distribution that approximates the BSDF is important for efficient rendering. Hery and Ramamoorthi (2012) showed how to sample the first term of the Marschner et al. model, and d'Eon et al. (2013) showed how to sample all terms of their improved model. (See also Jakob (2012) for notes related to sampling their  $M_p$  term in a numerically stable way.)

d'Eon et al. (2014) performed extensive Monte Carlo simulations of scattering from dielectric cylinders with explicitly modeled scales and glossy scattering at the boundary based on a Beckmann microfacet distribution. They showed that separable models didn't model all of the observed effects and that in particular that the specular term modeled by  $M_p$  varies over the surface of the cylinder and also depends on  $\phi$ . They developed a non-separable scattering model, where both  $\alpha$  and  $\beta_m$  varied as a function of h, and showed that it fit their simulations very accurately.

All of the scattering models we have described so far have been BCSDFs—they represent the overall scattering across the entire width of the hair in a single model. Such "far field" models assume that both the viewer is far away and that incident illumination is uniform across the hair's width. In practice, both of these assumptions are invalid if one is using path tracing to model multiple scattering inside hair. Two recent models have considered scattering at a single point along the hair's width, making them more suitable for accurately modeling "near field" scattering.

Yan et al. (2015) generalized d'Eon et al.'s model to account for scattering in the medulla, modeling a scattering cylinder in the interior of fur. They validated their model with a variety of measurements of actual animal fur, and showed how previous hair scattering models didn't match measured fur scattering well. Their model didn't integrate across the hair width.

Chiang et al. (2016) showed a number of comparisons that eliminating the integral over width from d'Eon et al.'s model works well in practice, and that the sampling rates necessary for path tracing also worked well to integrate scattering over the curve width, giving a much more efficient implementation. They also developed the perceptually-uniform parameterization of  $\beta_m$  and  $\beta_n$  that we have

implemented here as well as the inverse mapping from reflectance to  $\sigma_a$  used in our HairBSDF::SigmaAFromReflectance() method.

## **1.9 BIBLIOGRAPHY**

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## 1.11 EXERCISES

1.1 Marschner et al. (2003) note that human hair actually has an elliptical cross section that causes glints in human hair due to caustics. Extend the implementation here to handle this case. One issue that you'll need to address is that the  $\partial p/\partial v$  returned by Curve::Intersect() is always

perpendicular to the incident ray, which leads to different orientations of the azimuthal coordinate system. This isn't an issue for the model we have implemented here, since it operates only on the difference between angles  $\phi$  in the hair coordinate system. For elliptical hairs, a consistent azimuthal coordinate system is necessary.

- 1.2 Ogaki et al. (2010) created an explicit geometric model of the cuticle surface and then shot a large number of rays at it for each of a set of discrete outgoing directions, modeling scattering at the boundary with a microfacet model and modeling absorption and scattering in the hair interior. They then created a tabular representation of the resulting scattering distribution and used it for rendering. Implement this approach and compare the result to the model here.
- 1.3 As discussed in "A Note on Reciprocity", the model implemented in this document doesn't obey reciprocity. Investigate this issue and derive an improved model that does.
- 1.4 Read Yan et al.'s paper on fur scattering (2015) and implement their model, which accounts from scattering in the medulla in fur. Render images that show the difference from accounting for this in comparison to the current implementation. You may want to also see Section 4.3 of Chiang et al. (2016), which discusses extensions for modeling the undercoat (which is shorter and curlier hair underneath the top level) and a more ad-hoc approach to account for the influence of scattering from the medulla.
- 1.5 Read the paper by d'Eon et al. (2014) on a non-separable hair scattering model and implement their approach in pbrt. Render images that show the difference between their approach and the current implementation.

## 1.12 REVISION HISTORY

October 16, 2016: original version posted. November 5, 2016: very minor typos fixed.

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